# NASA TECHNICAL MEMORANDUM

NASA TM X-53249 April 23, 1965

IASA TM X-53249

GPO PRICE \$	
CSFTI PRICE(S) \$	
Hard copy (HC) _	1.00
Microfiche (MF)	<del></del>

# NEAR-OPTIMUM GUIDANCE - AN ANALYSIS OF FUEL PENALTY

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by Lyle R. Dickey Aero-Astrodynamics Laboratory

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George C. Marshall Space Flight Center, Huntsville, Alabama

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Ву

Lyle R. Dickey

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Huntsville, Alabama

#### ABSTRACT

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An explicit solution of the linearized equations of motion is used to obtain the deviations in end conditions,  $\triangle r$  and  $\triangle \theta$ , and the additional burning time,  $\triangle t$ , expressed as integrals of functions of thrust angle deviations,  $\triangle x$ . Under the constraint that  $\triangle r = \triangle \theta = 0$ , the calculus of variations is applied to minimize  $\triangle t$ . The resulting Euler-Lagrange equation evaluated at  $\triangle x = 0$  gives a simple relationship which is used to show that  $\triangle t$  is a second order function of  $\triangle x$ . This function is evaluated numerically for the second stage of an early SA-6 design. Results show that, if the mission is accomplished, the thrust angle may differ by as much as 2 degrees throughout the second stage with a propellant penalty of only 34 pounds. It follows that the capability of meeting the mission is the prime requisite of a guidance function and optimality is then only a second order consideration.

## NASA, GEORGE C. MARSHALL SPACE FLIGHT CENTER

Technical Memorandum X-53249

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Ву

Lyle R. Dickey

TECHNICAL AND SCIENTIFIC STAFF
AERO-ASTRODYNAMICS LABORATORY
RESEARCH AND DEVELOPMENT OPERATIONS

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## DEFINITION OF SYMBOLS

Symbol Symbol	<u>Definition</u>
t	time measured on a standard trajectory
to	second stage ignition time on the standard trajectory
t <sub>n</sub>	second stage cutoff time on the standard trajectory
w	propellant flow rate
∆r	radius error at cutoff
∆t	deviation in second stage burning time
$\triangle t_{o}$	deviation in second stage ignition time from standard
Δ₩	deviation in propellant consumption from standard
Δθ	error in angle between the velocity vector and the position vector at cutoff
ΔΧ	$\chi(t + \Delta t_0) - \chi_s(t)$
χ	thrust angle
Χ <sub>s</sub>	thrust angle on the standard trajectory

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#### SUMMARY

An explicit solution of the linearized equations of motion is used to obtain the deviations in end conditions,  $\Delta r$  and  $\Delta \theta$ , and the additional burning time,  $\Delta t$ , expressed as integrals of functions of thrust angle deviations,  $\Delta X$ . Under the constraint that  $\Delta r = \Delta \theta = 0$ , the calculus of variations is applied to minimize  $\Delta t$ . The resulting Euler-Lagrange equation evaluated at  $\Delta X = 0$  gives a simple relationship which is used to show that  $\Delta t$  is a second order function of  $\Delta X$ . This function is evaluated numerically for the second stage of an early SA-6 design. Results show that, if the mission is accomplished, the thrust angle may differ by as much as 2 degrees throughout the second stage with a propellant penalty of only 34 pounds. It follows that the capability of meeting the mission is the prime requisite of a guidance function and optimality is then only a second order consideration.

#### I. INTRODUCTION

The task of constructing a guidance function to determine the thrust angle, X, that meets a given mission under the constraint of minimum propellant consumption can be considerably simplified by first investigating the penalty for non-optimality. It is assumed that minimum propellant consumption is equivalent to minimum burning time, and it will be shown that any additional burning time required is a second order function of thrust angle deviations from optimum provided the mission is met. A numerical example is included which covers the second stage of a 100 n.m. orbital mission for an early SA-6 design. The results show that any guidance function which meets the required end conditions is acceptable even though it differs from the calculus of variations solution by as much as one or two degrees throughout the entire second stage.

#### II. EULER-LAGRANGE EQUATIONS

It is assumed that the calculus of variations solution is known for a given trajectory. If the thrust angle, X, deviates slightly from optimum, it has been shown in Reference 1 that the errors in end conditions  $\triangle \mathbf{r}$  and  $\triangle \theta$  and the additional burning time  $\triangle \mathbf{t}$  can be adequately expressed as follows:

$$\Delta \mathbf{r} = \int_{0}^{t_{n}} f(\Delta \mathbf{x}, t) dt, \qquad (1)$$

$$\triangle \theta = \int_{t_0}^{t_n} g(\triangle X, t) dt, \qquad (2)$$

$$\triangle t = \int_{t_0}^{t_n} h(\triangle x, t) dt, \qquad (3)$$

where

$$f(\Delta X, t) = f_1(t) \Delta X + f_2(t) \Delta X^2 + \dots,$$
 (4)

$$g(\Delta X,t) = g_1(t) \Delta X + g_2(t) \Delta X^2 + \dots , \qquad (5)$$

$$h(\Delta X, t) = h_1(t) \Delta X + h_2(t) \Delta X^2 + \dots$$
 (6)

Minimizing equation (3) under the constraint that equations (1) and (2) satisfy prescribed conditions leads to the Euler-Lagrange equation [2, p. 51]

$$\frac{\partial f^*}{\partial x} - \frac{d}{dt} \frac{\partial f^*}{\partial x} = 0 \tag{7}$$

where

$$f^* = h + \lambda_1 f + \lambda_2 g. \tag{8}$$

Since  $f^*$  is independent of  $\triangle x$ , equation (7) degenerates to the following:

$$\frac{\partial \mathbf{f}^*}{\partial \Delta X} = \frac{\partial \mathbf{h}}{\partial \Delta X} + \lambda_1 \frac{\partial \mathbf{f}}{\partial \Delta X} + \lambda_2 \frac{\partial \mathbf{g}}{\partial \Delta X} = 0,$$

which yields the following necessary condition:

$$(h_1 + \lambda_1 f_1 + \lambda_2 g_1) + 2(h_2 + \lambda_1 f_2 + \lambda_2 g_2) \triangle X + \dots = 0.$$
 (9)

Since the standard was assumed to be a calculus of variations solution, equation (9) is satisfied for  $\triangle X = 0$  and the following relationship is determined:

$$h_1 + \lambda_1 f_1 + \lambda_2 g_1 = 0.$$
 (10)

Table I shows the numerical values obtained for  $h_1$ ,  $f_1$ , and  $g_1$ . The following values of  $\lambda_1$  and  $\lambda_2$  were obtained by the method of least squares.

$$\lambda_1 = -.08734 \text{ sec/km}$$

$$\lambda_2 = .8589 \text{ sec/deg.}$$

These values satisfy equation (10) within the numerical accuracy with which  $h_1,\ f_1,\ and\ g_1$  were determined. This is illustrated in Table I where the quantity  $h_1+\lambda_1f_1+\lambda_2g_1$  is shown.

#### III. SECOND ORDER EFFECT

From equation (10), it follows that

$$h_1 \triangle X + \lambda_1 f_1 \triangle X + \lambda_2 g_1 \triangle X = 0$$

TABLE I

t (sec)	$(10^{-2} \text{ km/deg sec})$	(10 <sup>-2</sup> /sec)	(10 <sup>-2</sup> /deg)	$h_1 + \lambda_1 f_1 + \lambda_2 g_1$ (10 <sup>-2</sup> /deg)
160	<b>-3.6</b> 80	.0454	3610	0006
200	-3.738	.0536	3730	<b></b> 0005
240	-3.772	.0628	3836	0002
280	-3.774	.07 <b>3</b> 2	3927	0002
<b>3</b> 20	-3.738	.0850	3996	0001
<b>36</b> 0	-3.654	.0986	4038	.0000
400	-3.507	.1144	4044	.0002
440	-3.280	.1330	4005	.0002
480	-2.944	.1554	3903	.0003
520	-2.458	.1828	<b></b> 3715	.0002
560	-1.758	.2177	3405	.0000
600	727	.2643	2907	0002

 $\lambda_1 = -.08734 \text{ sec/km}$ 

 $\lambda_2$  = .8589 sec/deg.

and

$$\int_{t_{0}}^{t_{n}} h_{1} \triangle X dt = -\lambda_{1} \int_{t_{0}}^{t_{n}} f_{1} \triangle X dt - \lambda_{2} \int_{t_{0}}^{t_{n}} g_{1} \triangle X dt.$$

$$(11)$$

From equations (4) and (5) together with equations (1) and (2), the following relationships are determined:

$$- \lambda_1 \int_{t_0}^{t_n} f_1 \Delta x dt = -\lambda_1 \Delta r + \lambda_1 \int_{t_0}^{t_n} [f_2 \Delta x^2 + \dots] dt,$$

and

$$- \lambda_2 \int_{t_0}^{t_n} g_1 \Delta x dt = - \lambda_2 \Delta \theta + \lambda_2 \int_{t_0}^{t_n} [g_2 \Delta x^2 + \dots] dt.$$

Substituting these values into equation (11) yields

$$\int_{t_{0}}^{t_{n}} h_{1}\Delta x \ dt = - \lambda_{1}\Delta r - \lambda_{2}\Delta \theta + \int_{t_{0}}^{t_{n}} [(\lambda_{1}f_{2} + \lambda_{2}g_{2}) \ \Delta X^{2} + \dots] \ dt.$$

This, together with equations (3) and (6), gives the following expression for  $\triangle t$ .

$$\Delta t = -\lambda_{\underline{x}} \Delta \mathbf{r} - \lambda_{\underline{z}} \Delta \theta + \int_{0}^{t_{n}} [(h_{\underline{z}} + \lambda_{\underline{1}} f_{\underline{z}} + \lambda_{\underline{2}} g_{\underline{z}}) \Delta X^{\underline{z}} + \dots] dt.$$
 (12)

If the required end conditions are met,  $\triangle r = \triangle \theta = 0$ . Then,

$$\Delta t = \int_{0}^{t_{n}} [(h_{2} + \lambda_{1}f_{2} + \lambda_{2}g_{2}) \Delta X^{2} + \dots] dt,$$
 (13)

and the additional burning time  $\triangle t$  resulting from non-optimum X is a second order function of  $\triangle X$ . The sufficient condition that t is a local minimum is that  $h_2 + \lambda_1 f_2 + \lambda_2 g_2 > 0$ ,  $t_0 \le t \le t_n$ .

#### IV, NUMERICAL RESULTS

Table II shows the numerical values obtained for  $h_2$ ,  $f_2$ ,  $g_2$  and  $(h_2 + \lambda_1 f_2 + \lambda_2 g_2)$ . In addition, the quantity  $h_2^*(t_1) = (h_2 + \lambda_1 f_2 + \lambda_2 g_2)$   $\Delta t_1$  is shown in Table III. This quantity can be used to obtain a good approximation to the integral of the second ordered term in equation (13) by the following summation.

$$\Delta t = \sum_{i=1}^{12} h_2^*(t_i) \Delta X^2(t_i). \tag{14}$$

This expression, in turn, has the following upper bound:

$$\Delta t \leq \Delta X_{m}^{2} \sum_{i=1}^{12} h_{2}^{*}(t_{i}), \qquad (15)$$

which, for this example, is

$$\Delta t \leq .0409 \quad \Delta \chi_m^2,$$
 (16)

where  $\triangle X_m$  is the maximum value of  $|\triangle X(t)|$ ,  $t_0 \le t \le t_n$ .

The propellant loss  $\triangle W$  can be expressed as follows:

$$\wedge W = \dot{W} \wedge t$$
.

For this example,  $\dot{W}$  = 208 lbs/sec, which yields the following bound:

$$\triangle W \leq 8.51 \triangle \chi_{m}^{2} \text{ (lbs/deg}^{2)}.$$
 (17)

TABLE II

t	$\mathtt{f}_{\mathcal{Z}}$	<b>g</b> <sub>2</sub>	$\mathfrak{h}_{2}$	$(h_2 + \lambda_1 f_2 + \lambda_2 g_2)$
(sec)	$(10^{-3}$ km/deg $^2$ sec)	(10 <sup>-5</sup> /deg sec)	(10 <sup>-4</sup> /deg <sup>2</sup> )	(10 <sup>-4</sup> /deg <sup>2</sup> )
160	3831	.8899	.2168	.6278
200	3516	.8798	.2606	.6433
240	3194	.8663	.3094	.6628
280	2863	.8482	.3641	.6870
320	2524	.8239	.4260	.7172
360	2175	.7915	.4969	.7548
400	1819	.7483	.5789	.8021
440	1458	.6907	.6755	.8622
480	1096	.6131	.7915	.9399
520	0742	.5074	.9344	1.0428
560	0408	.3602	1.1167	1.1832
600	0118	.1485	1.3604	1.3834

 $\lambda_1 = -.08734$ 

 $\lambda_2$  = .8589 sec/deg

# TABLE III

t (sec <b>)</b>	∆t i (sec)	h <sup>*</sup> <sub>2</sub> (t <sub>i</sub> ) (10 <sup>-2</sup> sec/deg <sup>2</sup> )
160	<b>33.</b> 19	.2084
200	40	.2573
240	40	.2651
280	40	. 2748
<b>3</b> 20	40	.2869
360	40	.3019
400	40	.3208
440	40	. 3449
480	40	. 3760
520	40	.4171
560	40	.4733
600	40.68	.5628
		$\sum_{i=1}^{12} h_2^*(t_i) = 4.0893 \times 10^{-2} \text{ sec/deg}^2$

#### V. CONCLUSIONS

Equation (17) shows that the propellant penalty from non-optimum guidance is less than 34 pounds even if the thrust angle deviates from optimum by as much as 2 degrees throughout the second stage. This clearly indicates that, if a guidance function can be determined which meets the mission within satisfactory limits and predicts X within one or two degrees of the optimum value, there is very little to be gained by expending much effort to further reduce the prediction error in X.

#### REFERENCES

- Dickey, Lyle R., "Guidance Applications of Linear Analysis," NASA TM X-53166, George C. Marshall Space Flight Center, November 27, 1964, unclassified.
- 2. Weinstock, Robert, "Calculus of Variations," McGraw-Hill, New York, 1952.

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